



**PARTICLE ACCELERATION IN
SUPERNOVA REMNANTS**

AND

THE ORIGIN OF COSMIC RAYS

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AGILE MEETING – April 16-17, 2012

WHICH QUESTIONS TO ANSWER TO?

ARE CRs ACCELERATED IN SNRs WITH AN EFFICIENCY AND SPECTRAL SHAPE COMPATIBLE WITH B/C AND WITH ANISOTROPY ?

ARE THE EXISTING GAMMA RAY OBSERVATIONS (AND THE WHOLE MULTI-■ SHEBANG) UNDERSTOOD IN TERMS OF KNOWN PHYSICS OF PARTICLE ACCELERATION?

DO WE SEE (and CAN WE SEE) EVIDENCE FOR CR ACCELERATION UP TO THE KNEE (IN PROTONS)?

WHAT IS THE CHEMICAL COMPOSITION AT THE SOURCE AND AT EARTH?

WHERE DO GALACTIC CRs END?

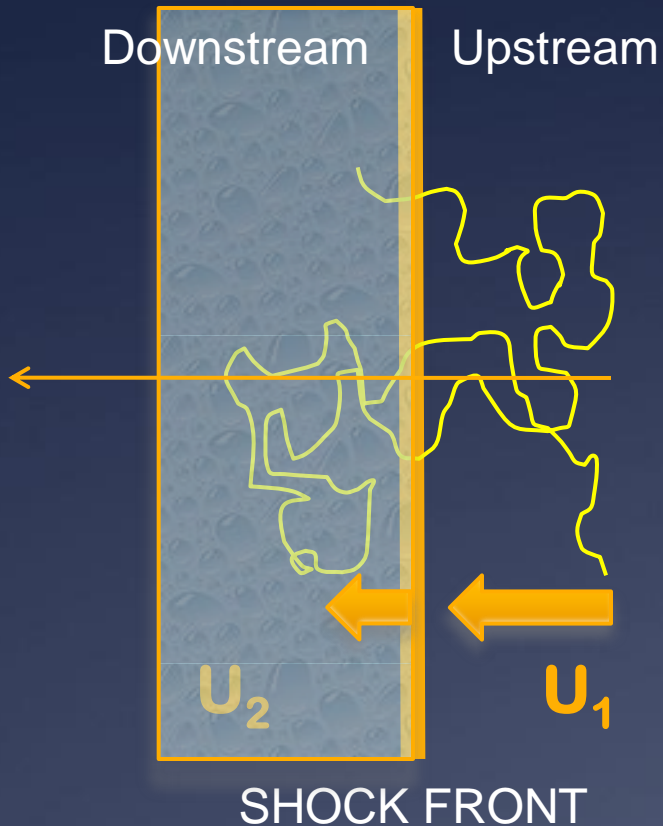
THEORY OF CR ACCELERATION IN SNRs

Diffusive Shock Acceleration

Berezhko, Ellison, Voelk; end of nineties +
PB 2002, 2004, Amato & PB 2005, 2006, Caprioli +, 2007-2011, Morlino + 2007-2012

First order FERMI ACCELERATION

Test particle theory



**DIFFUSION OF CHARGED PARTICLES
BACK AND FORTH THROUGH THE
SHOCK LEADS TO**

$$\frac{\Delta E}{E} = \frac{4}{3}(U_1 - U_2)$$

***PARTICLES ARE ACCELERATED TO A
POWER LAW SPECTRUM***

***THE SLOPE OF THE SPECTRUM ONLY
DEPENDS ON THE COMPRESSION***

NOT ON THE DIFFUSION COEFFICIENT

FOR STRONG SHOCKS: E^{-2}

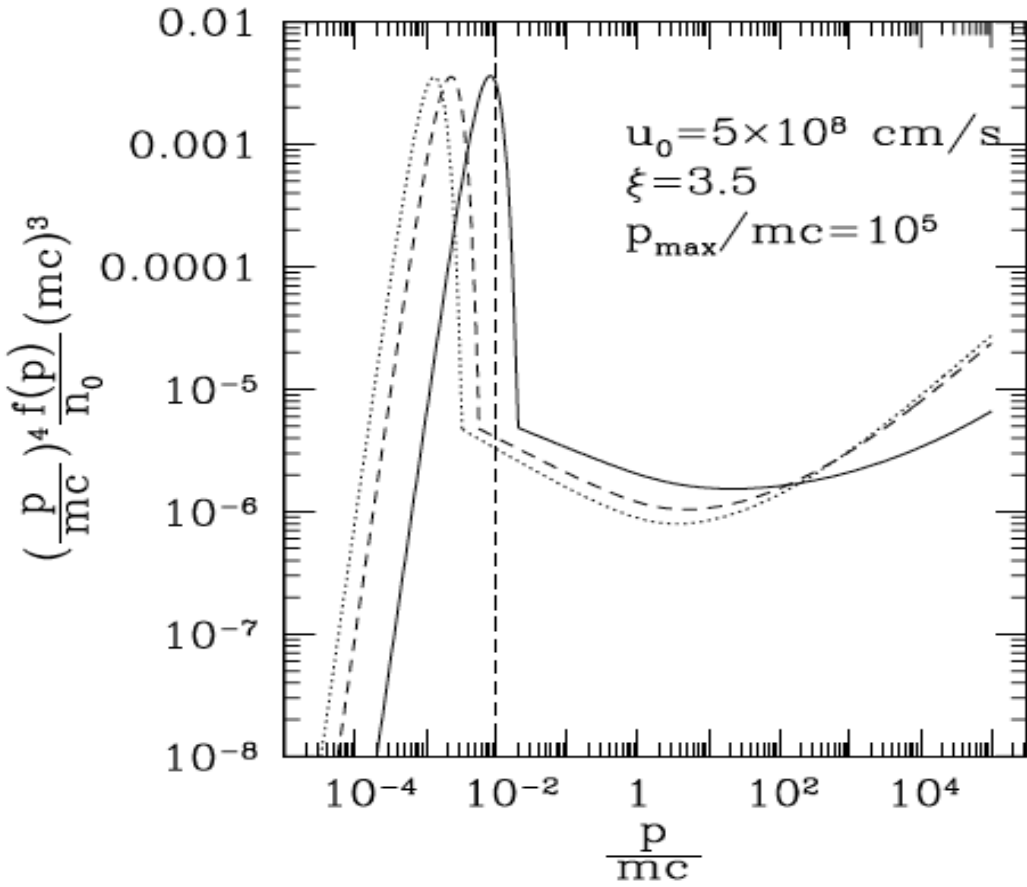
NON LINEAR THEORY

A theory of particle acceleration that allows one to describe:

- 1. Dynamical reaction of accelerated particles*
- 2. CR-induced B-field and their reaction*
- 3. Recipe for injection (self-regulation)*
- 4. Escape of particles (Cosmic Rays)*

DIFFUSIVE ACCELERATION AT COLLISIONLESS NEWTONIAN SHOCKS

non linear theory: BASIC PHYSICAL ASPECTS



**COMPRESSION FACTOR BECOMES
FUNCTION OF ENERGY**

**SPECTRA ARE NOT PERFECT
POWER LAWS (CONCAVE)**

**GAS BEHIND THE SHOCK IS
COOLER FOR EFFICIENT SHOCK
ACCELERATION**

SYSTEM SELF REGULATED

**EFFICIENT GROWTH OF B-FIELD
IF ACCELERATION EFFICIENT**

Particle Diffusion \leftrightarrow Wave Growth: STREAMING INSTABILITY

$$n_{CR} m v_D \rightarrow n_{CR} m V_A \Rightarrow \frac{dP_{CR}}{dt} = \frac{n_{CR} m (v_D - V_A)}{\tau}$$

$$\frac{dP_w}{dt} = \gamma_W \frac{\delta B^2}{8\pi} \frac{1}{V_A}$$

$$\gamma_W = \sqrt{2} \frac{n_{CR}}{n_{gas}} \frac{v_D - V_A}{V_A} \Omega_{cyc} \quad \text{GROWTH RATE}$$

In the ISM this is $\sim 10^{-3} \text{ yr}^{-1}$ but close to a shock front the growth can be much larger!!!

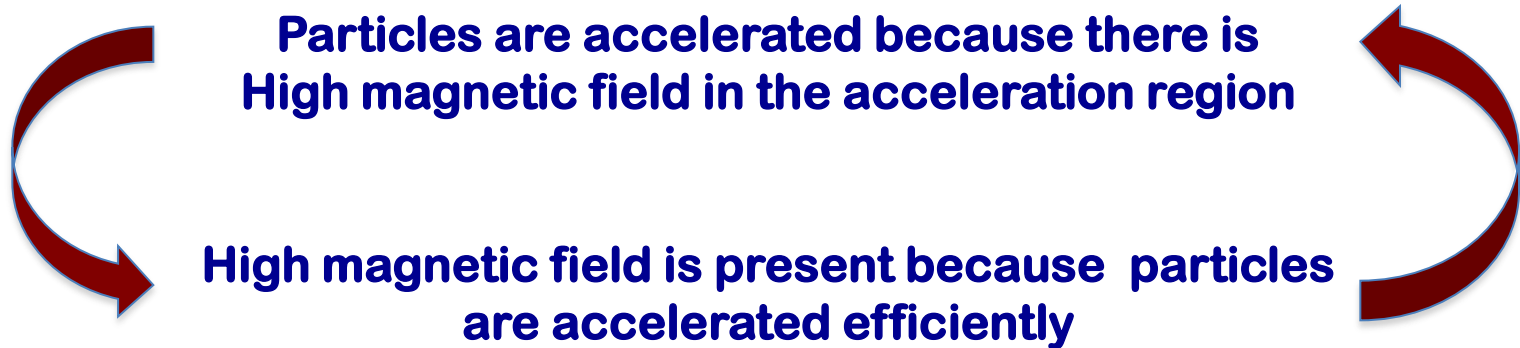
$$\gamma_W \simeq \sqrt{2} \xi_{CR} \left(\frac{V_s}{c} \right)^2 \frac{V_s}{V_A} \Omega_{cyc} \sim \mathcal{O}(\text{seconds}^{-1})$$

\underline{B} IS AMPLIFIED BY PARTICLES

MAGNETIC FIELD AMPLIFICATION

MAGNETIC FIELD AMPLIFICATION

SMALL PERTURBATIONS IN THE LOCAL B-FIELD CAN BE AMPLIFIED BY THE SUPER-ALFVENIC STREAMING OF THE ACCELERATED PARTICLES



Without this non-linear process, no acceleration of CR to High energies (and especially not to the knee!)

X-ray rims and B-field amplification

TYPICAL THICKNESS OF FILAMENTS: $\sim 10^{-2}$ pc

The synchrotron limited thickness is:

$$\Delta x \approx \sqrt{D(E_{max})\tau_{loss}(E_{max})} \approx 0.04 B_{100}^{-3/2} \text{ pc}$$


$$B \approx 100 \mu\text{Gauss}$$

$$E_{max} \approx 10 B_{100}^{-1/2} u_8 \text{ TeV}$$

$$\nu_{max} \approx 0.2 u_8^2 \text{ keV}$$

In some cases the strong fields are confirmed by time variability of X-rays
Uchiyama & Aharonian, 2007

TROUBLE WITH SPECTRA

THE SPECTRA OF ACCELERATED PARTICLES ARE IN GENERAL CONCAVE AND FLATTER THAN E^{-2} AT HIGH ENERGY

THE MAXIMUM ENERGY WITH B-FIELD AMPLIFICATION MAY REACH UP TO $\sim 10^{15}$ eV FOR PROTONS (Z TIMES HIGHER FOR NUCLEI)

THESE SPECTRA SHOULD REFLECT IN THE GAMMA RAY SPECTRA (IF DUE TO PP SCATTERING) AND NEUTRINO SPECTRA

BUT THE OBSERVED SPECTRA OF GAMMAS ARE TYPICALLY $\sim E^{-2.3}$

OFTEN INCOMPATIBLE WITH LEPTONIC MODELS! BUT ALSO NOT COMPATIBLE WITH THE SIMPLEST PREDICTION OF NLDSA

THE EFFECT OF THE WAVE SPEED?

One should remember that the compression factor that counts in shock acceleration is not that of fluid velocity, but that of the scattering centers velocity

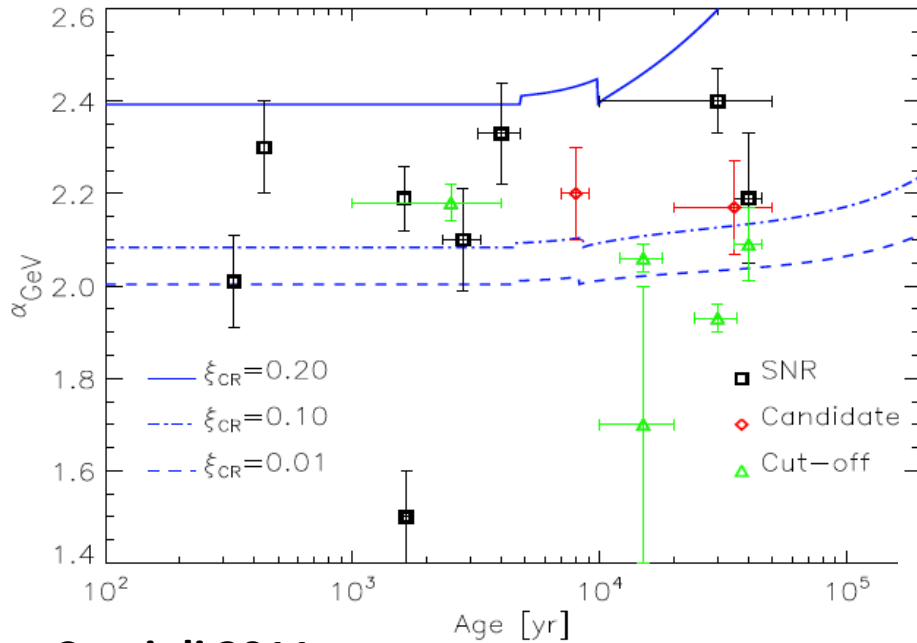
$$r = \frac{u_1}{u_2} \rightarrow \tilde{r} = \frac{u_1 + v_{A,1}}{u_2 + v_{A,2}}$$

When the magnetic field is amplified the Alfvén speed is not well defined and one may argue that it should be calculated in the amplified field (it depends on helicity!):

$$\tilde{r} = r \left(1 - \frac{1}{M_{A,1}} \right) = \frac{\gamma_{\text{eff}} + 1}{\gamma_{\text{eff}} - 1 + 2/M_s^2} \left[1 - \frac{\xi_{cr}(2 - \xi_{cr})}{2(1 - \xi_{cr})^{5/2}} \right] \quad \gamma_{\text{eff}} = \frac{1}{3} \frac{5 + 3\xi_{cr}}{1 + \xi_{cr}}$$

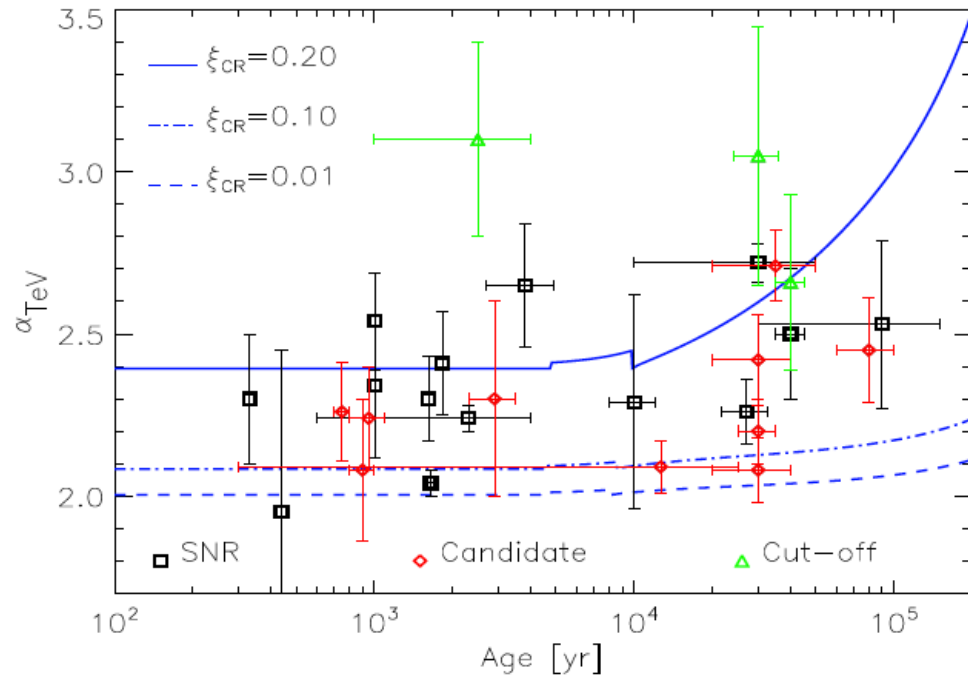
THIS EFFECT MAY LEAD TO STEEPER SPECTRA WHEN ACCELERATION IS EFFICIENT (BUT VERY MODEL DEPENDENT...could lead to opposite effect depending on wave helicity)

TROUBLE WITH SLOPES ?

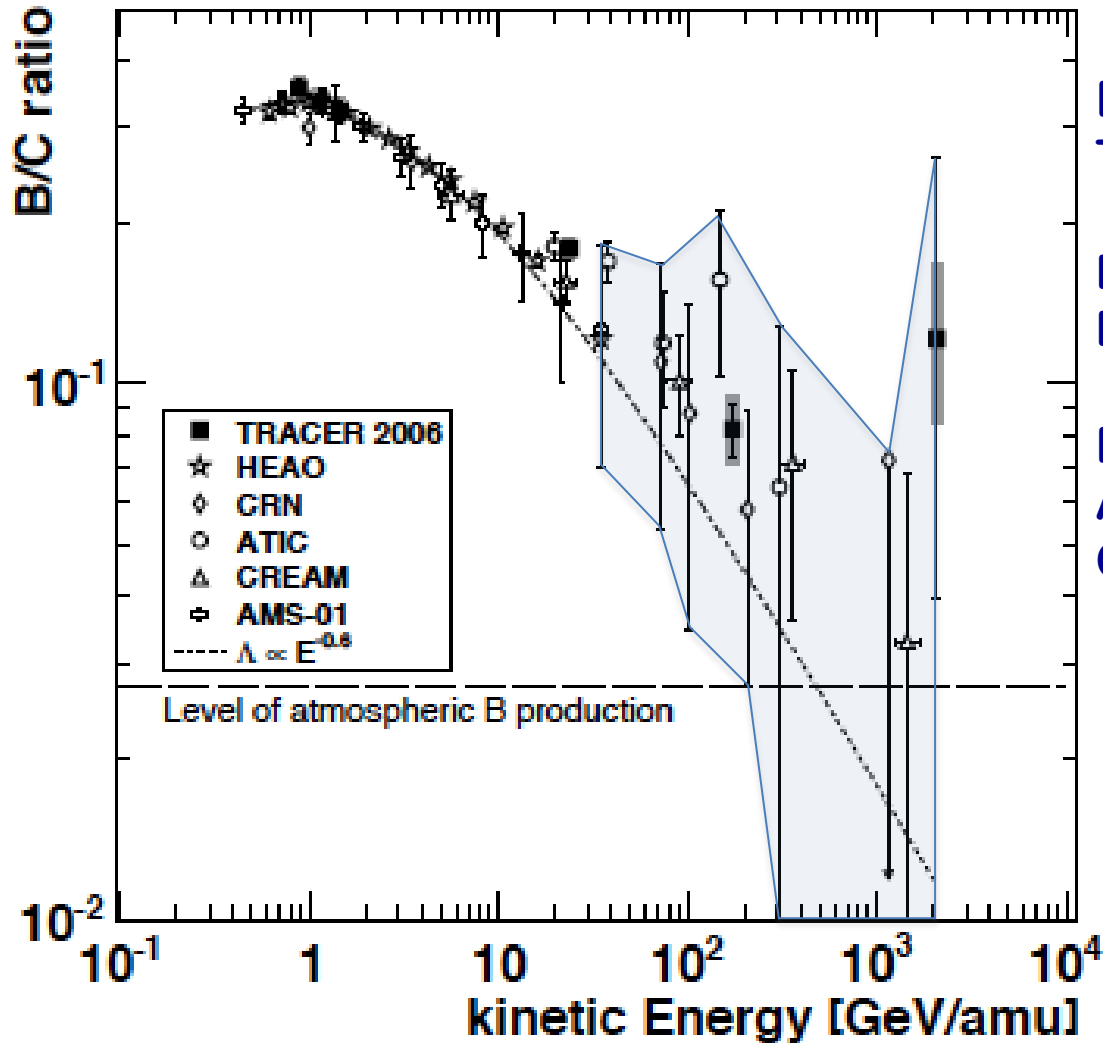


Caprioli 2011

VERY SURPRISING TO SEE THAT THE REQUIRED ACCELERATION EFFIC. ARE HIGH BUT THE SPECTRA ARE STEEP



IMPLICATIONS FOR CR PROPAGATION



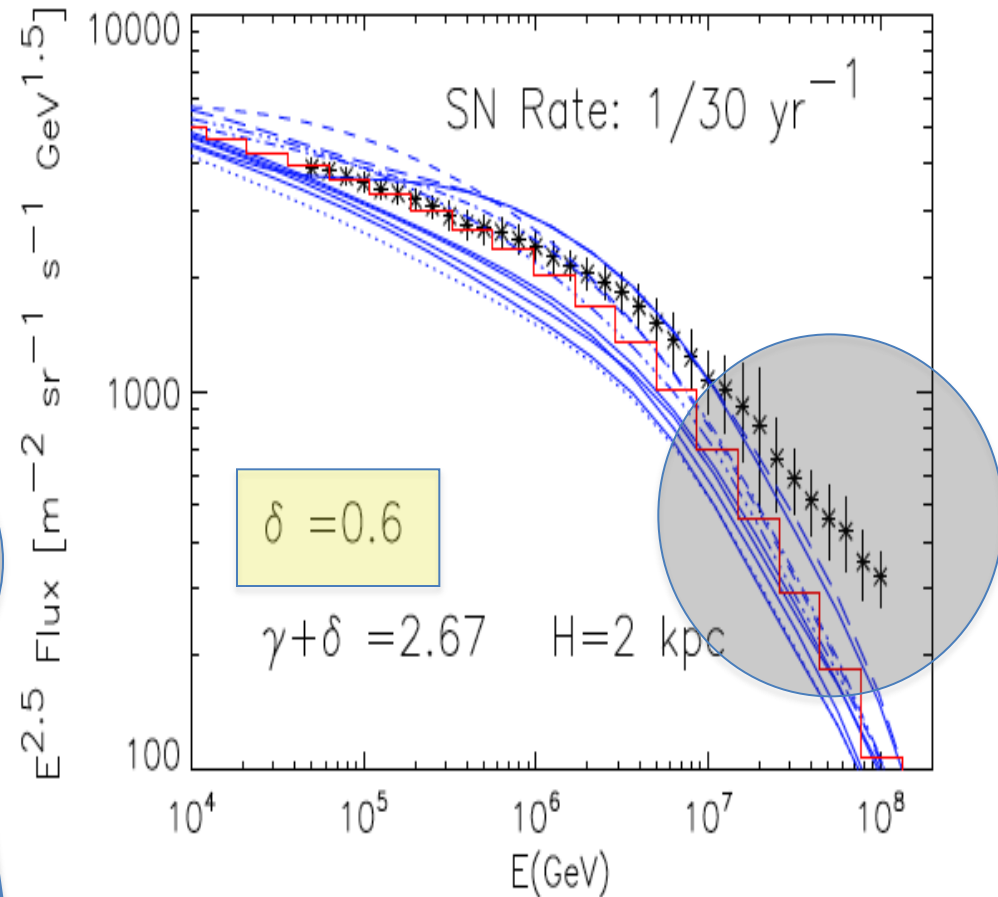
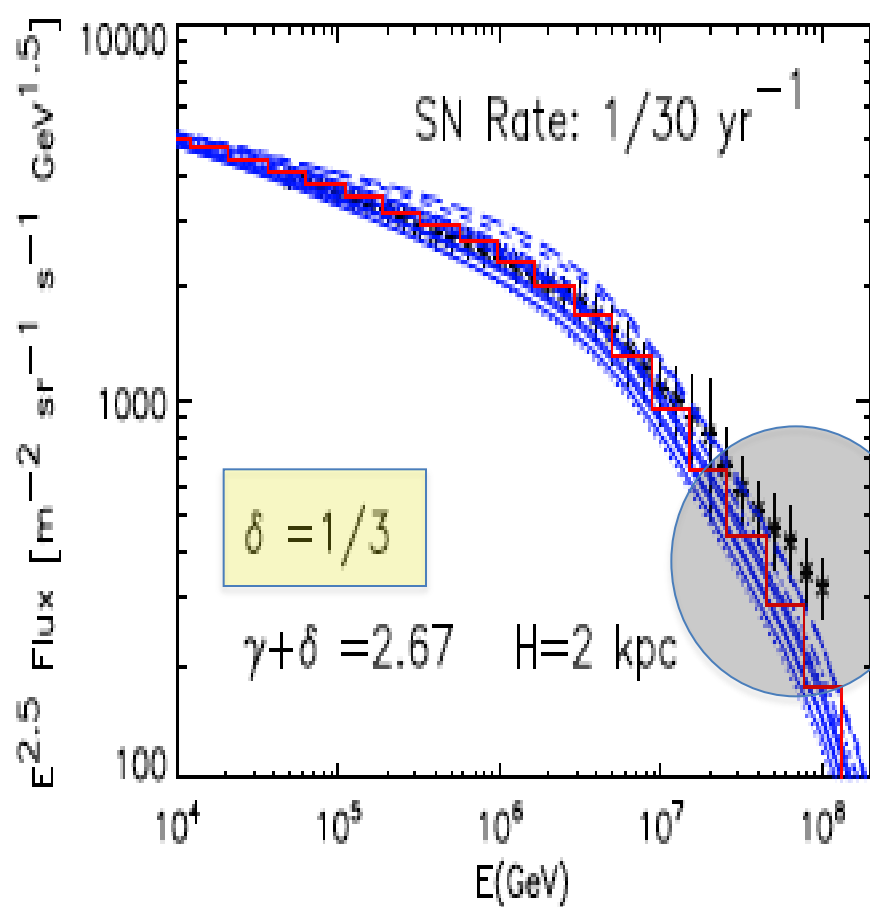
IN PRINCIPLE $B/C \sim 1/D(E)$ IN THE HIGH RIGIDITY REGIME

BUT UNCERTAINTIES ARE STILL LARGE

NOT EASY TO DISCRIMINATE AMONG DIFFERENT DIFFUSION COEFFICIENTS

CR spectra and SNRs

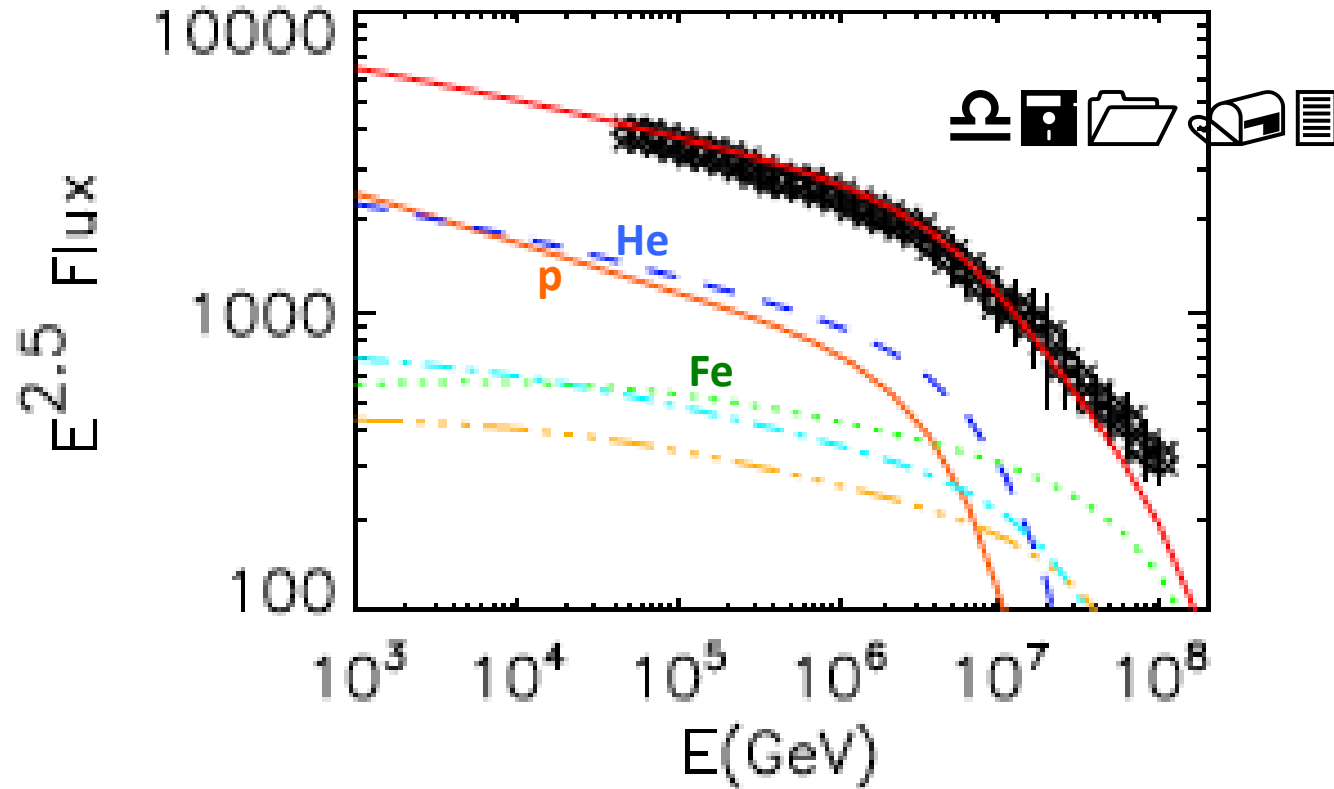
Blasi & Amato 2011



Deficit compensated
by extragalactic CRs

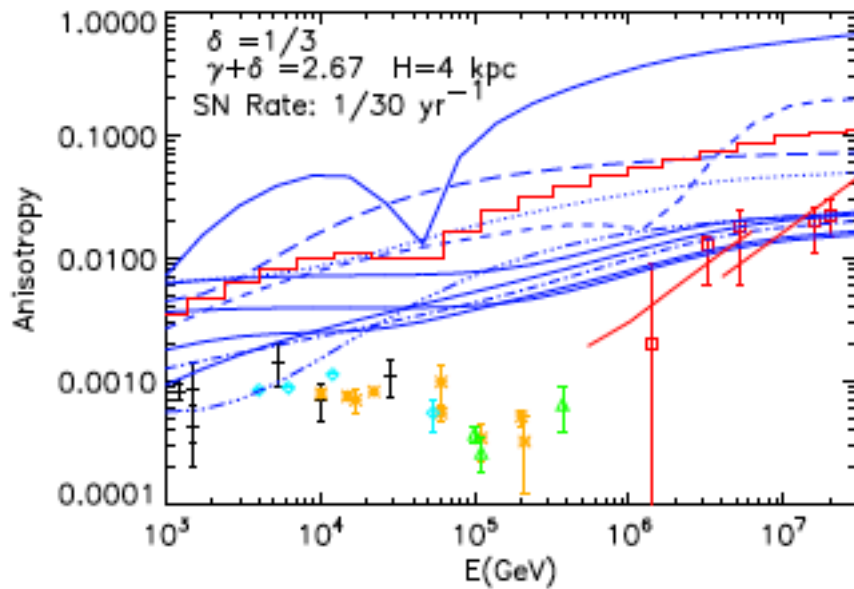
Hardening of Helium?

Blasi & Amato 2011



ONLY FOR $\Omega=1/3$ SPECTRUM OF He HARDER THAN SPECTRUM OF PROTONS AS A RESULT OF MODEST SPALLATION

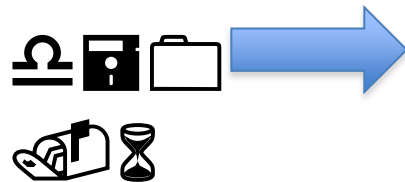
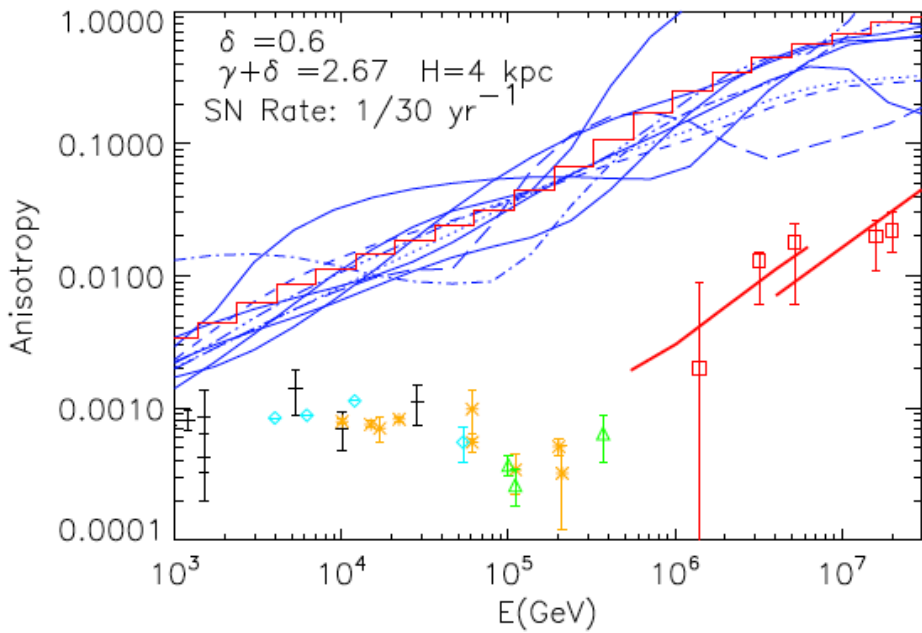
Large Scale CR Anisotropy



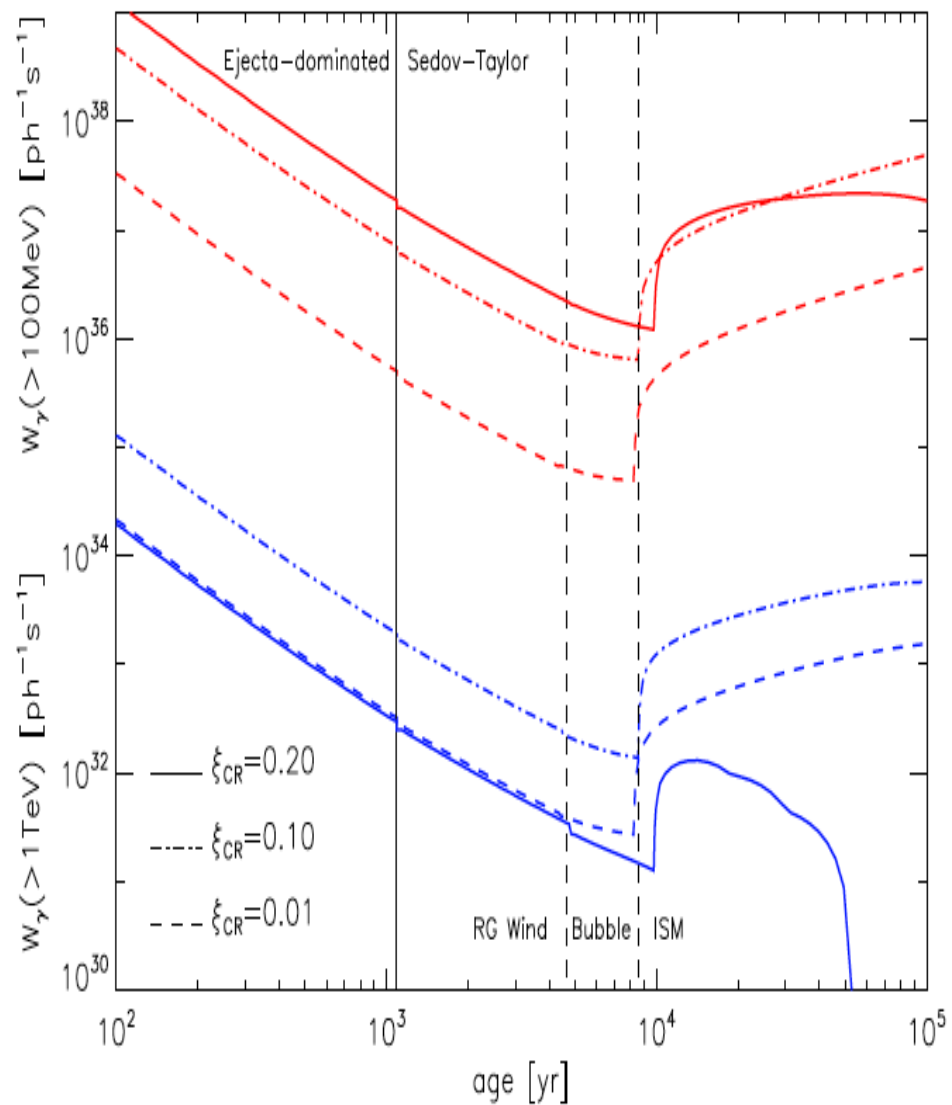
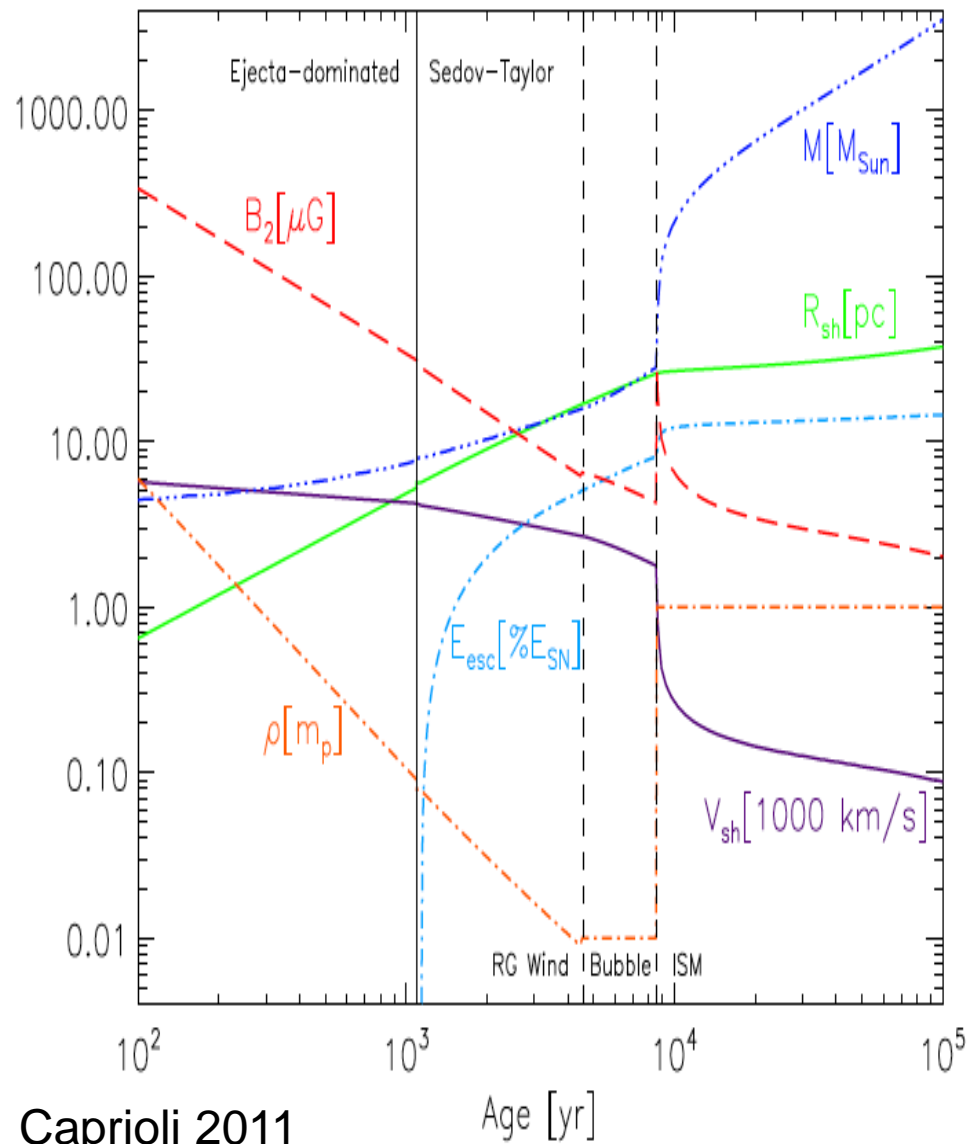
Naïve expectation:

$$\delta_A = \frac{3}{2^{3/2}} \frac{1}{\pi^{1/2}} \frac{D(E)}{Hc}$$

proportional to E^Ω



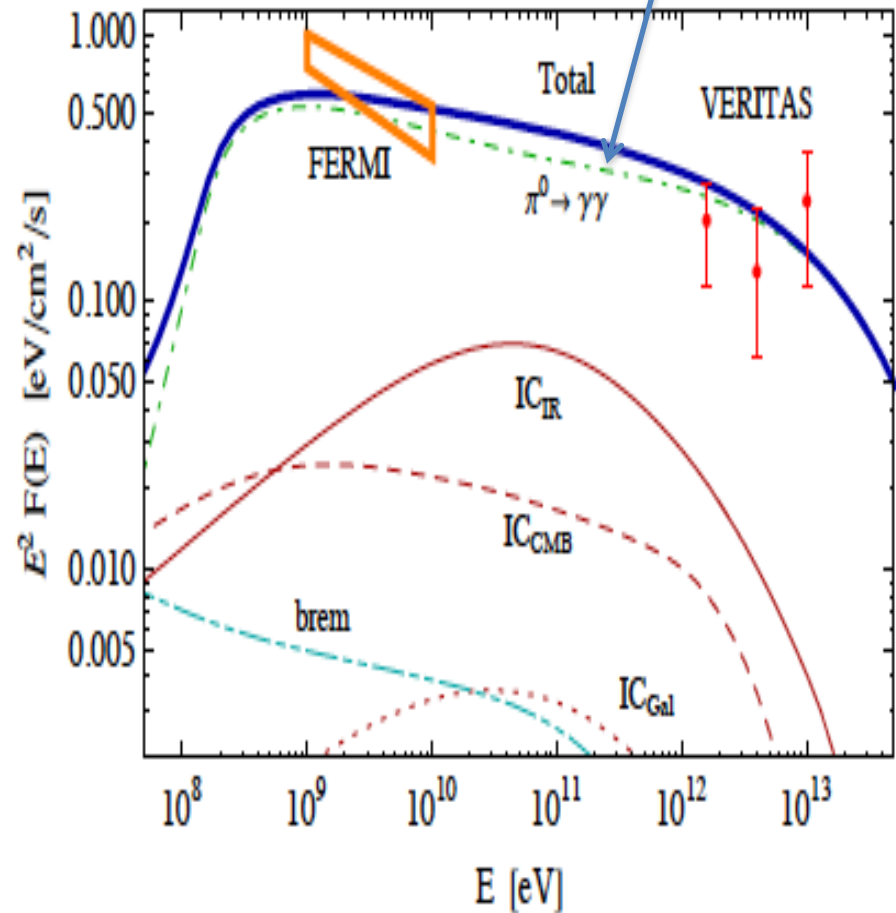
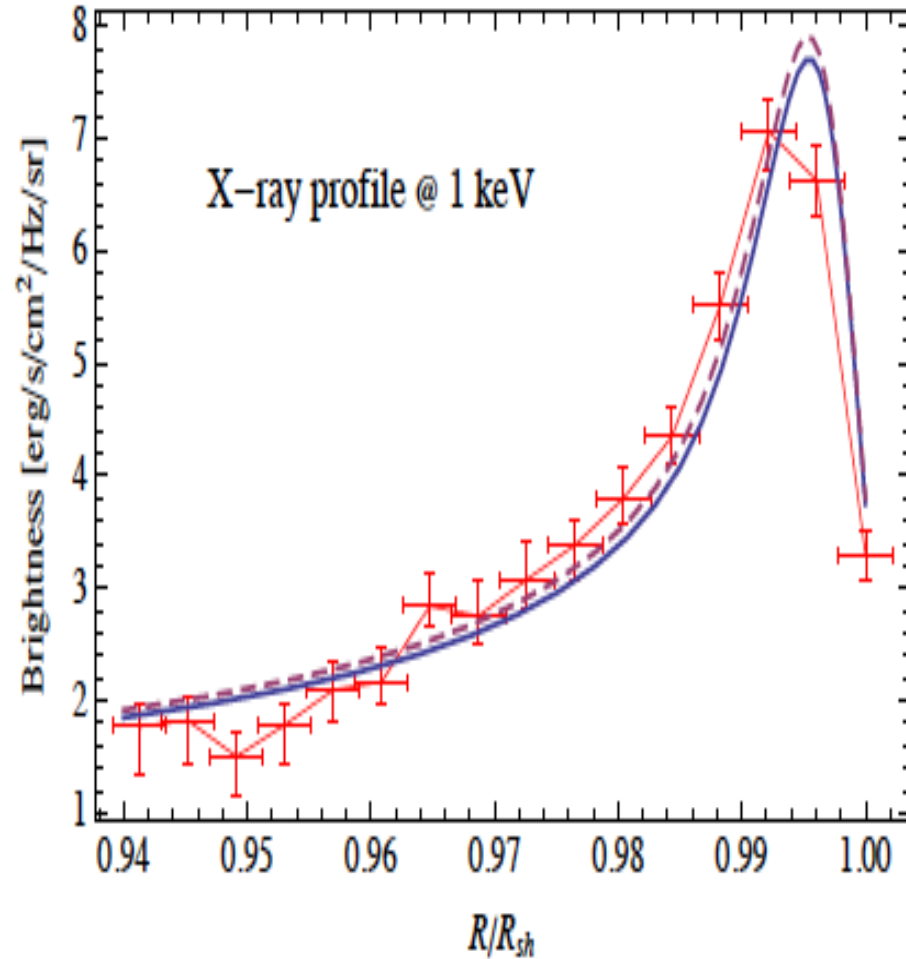
TEMPORAL EVOLUTION



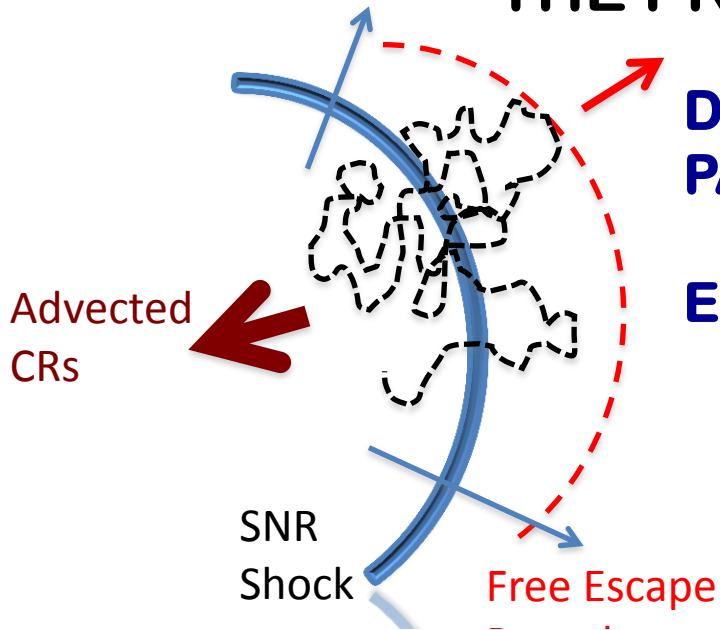
The prototypical case of Tycho

Morlino&Caprioli 2011

**STEEP SPECTRUM
BASICALLY IMPOSSIBLE TO
EXPLAIN WITH LEPTONS**



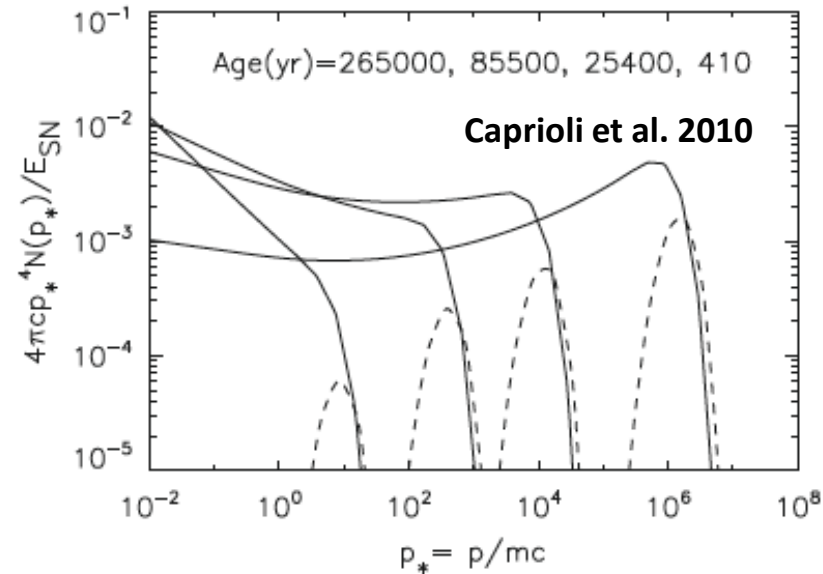
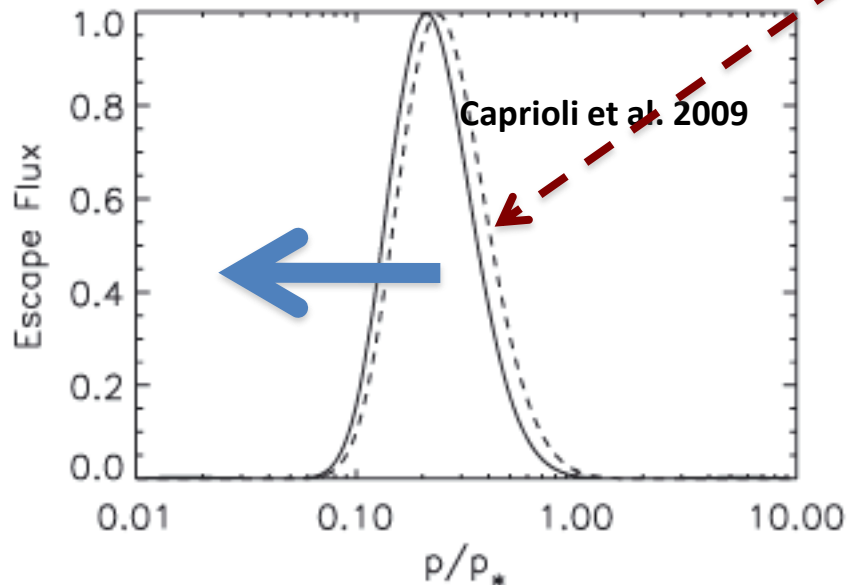
HOW DO ACCELERATED PARTICLES BECOME CRs? THE PROBLEM OF ESCAPE



**DURING THE EJECTA DOMINATED PHASE
PARTICLES ARE ADVECTED WITH PLASMA**

ESCAPE OCCURS DURING SEDOV PHASE

$$\Phi_{esc}(E, x) = D(E) \left(\frac{\partial f(E, x)}{\partial x} \right)_{x=x_{fe}}$$



CR ESCAPE AND CLOUDS

TWO SCENARIOS:

SNR SHOCK ENTERS THE MOLECULAR CLOUD

Collisionless shock only involves ions (very low density)

Ion-neutral density kills waves \rightarrow low E_{\max}

$$\text{Rate of damping } \Gamma_{ion-n} = \frac{1}{2} \langle n_H \sigma \rangle \approx 4.2 \times 10^{-9} T_4^{0.4} n_1 \text{ s}^{-1}$$

Remember: $D(E) = (1/3) c r_L / F(k)$ and here $F(k) \rightarrow 0$ in clouds

MOLECULAR CLOUD IS ILLUMINATED BY CR FROM SNR

The MC only acts as a target for pp. Gamma ray flux depends on

- Age of SNR
- Diffusion coefficient around the SNR (LIKELY SELF GENERATED IF CR DENSITY DOMINATED BY LOCAL SNR)
- Escape physics

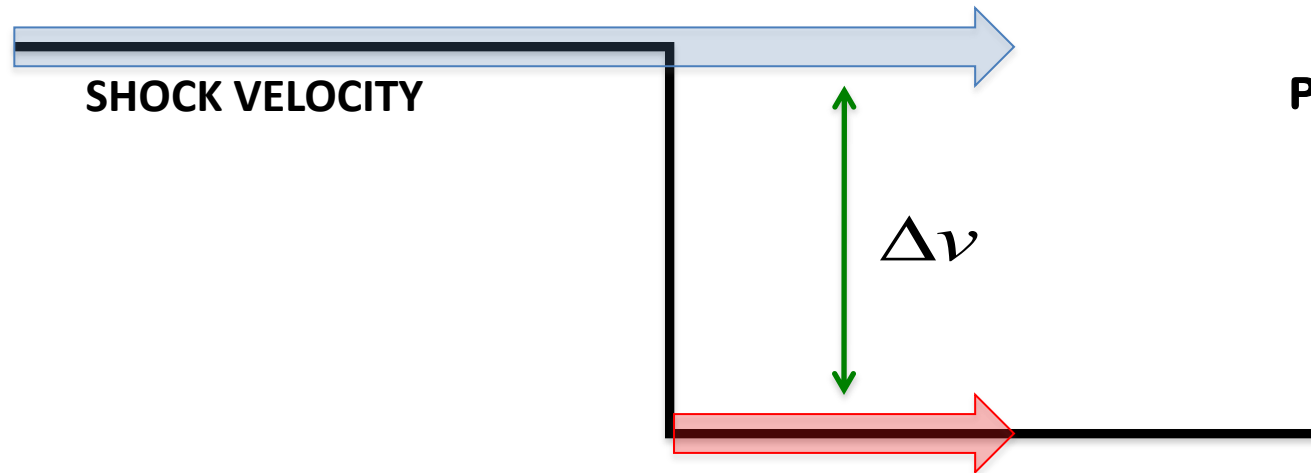


**SOME NEW WAYS TO MEASURE THE
CR ACCELERATION EFFICIENCY IN
SUPERNOVA REMNANTS**

***COLLISIONLESS SNR SHOCKS IN
PARTIALLY IONIZED MEDIA***

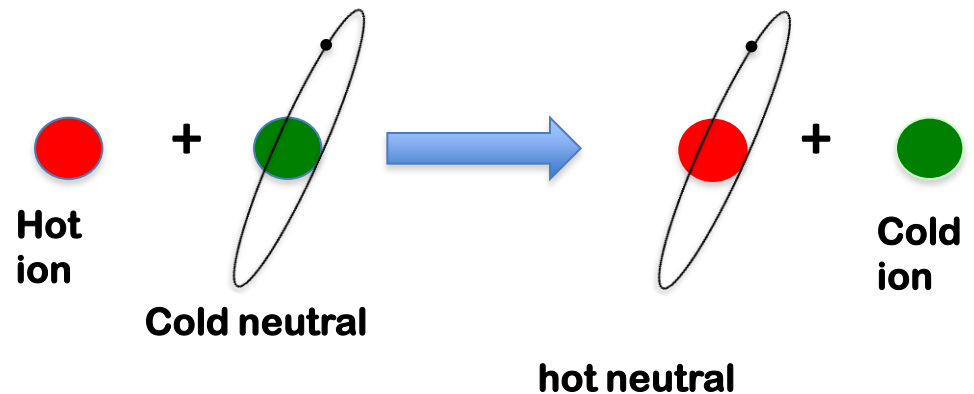
SUBTLE ASPECTS OF ACCELERATION AT A COLLISIONLESS SHOCK

NEUTRALS
AND IONS



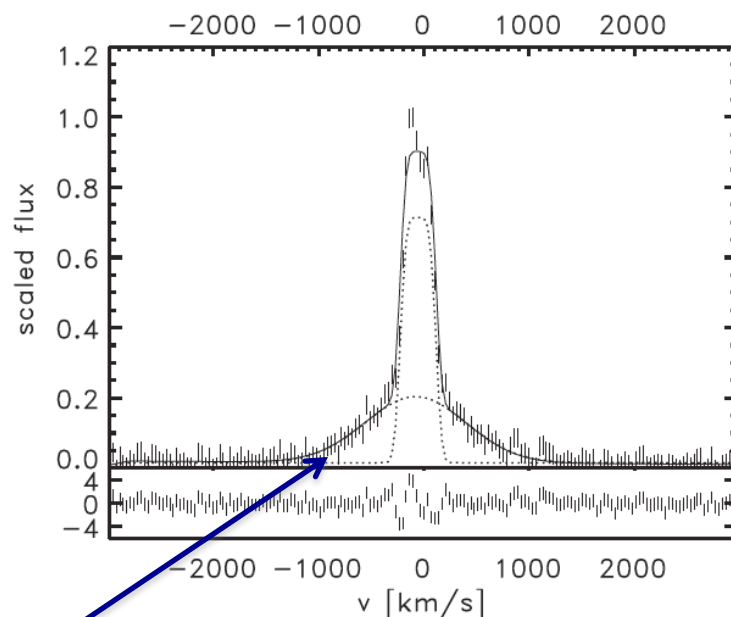
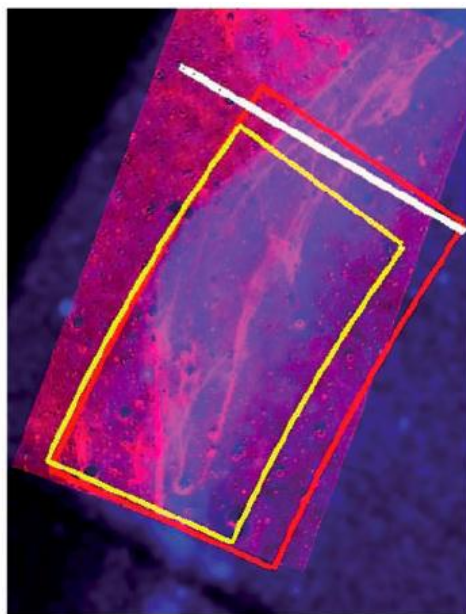
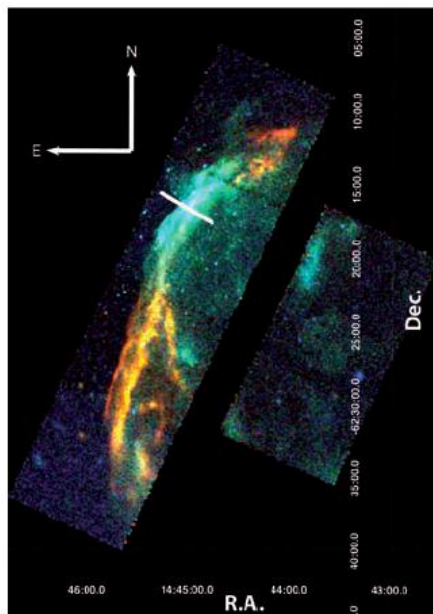
PB+, 2012

**CHARGE EXCHANGE → BROAD
BALMER LINE (NEUTRALS
THAT MADE CHARGE
EXCHANGE) REFLECTING
THE TEMPERATURE OF IONS...**



BUT THE LATTER AFFECTED BY EFFICIENT CR ACCELERATION

BROAD BALMER LINES NARROWER THAN FOR UNMODIFIED SHOCKS



Helder et al. 2009

$$W_{broad} = \sqrt{8 \ln 2 \frac{kT_2}{m}} \approx 1.02 v_{sh}$$

$$W_{broad} = 1100 \pm 63 \text{ km/s} \rightarrow T_2 = 2.3 \pm 0.3 \text{ keV}$$

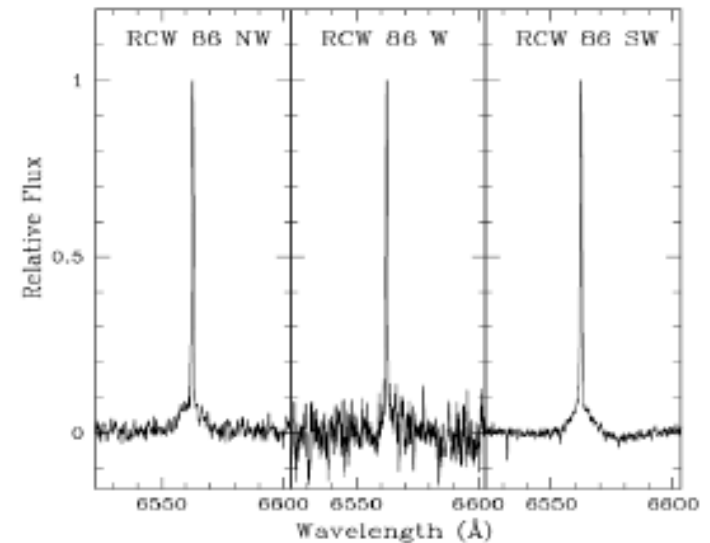
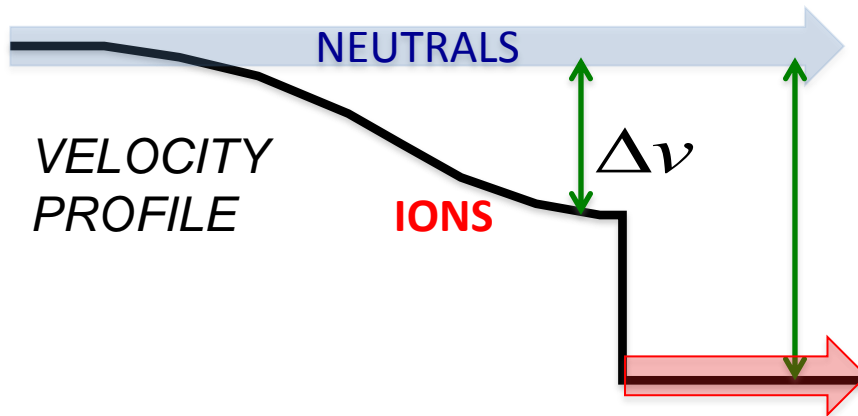
Shock speed from proper motion

$$v_{shock} = 6000 \pm 2800 \text{ km/s} \left(\frac{d}{2.5 \pm .5 \text{ kpc}} \right) \left(\frac{\dot{\theta}_{obs}}{0.5 \pm .2'' \text{ yr}^{-1}} \right) \rightarrow T_2 = \begin{matrix} 20-150 \text{ keV} (no equilibration) \\ 12-90 \text{ keV} (equilibration) \end{matrix}$$

INFERRED EFFICIENCY of CR ACCELERATION 50-60% !!! (BUT model dependent)

NARROW BALMER LINES BROADER THAN FOR UNMODIFIED SHOCKS

Sollerman et al. 2003



CHARGE EXCHANGE OCCURS NOW IN THE CR INDUCED PRECURSOR

$$W_{broad} = \sqrt{8 \ln 2 \frac{kT_0}{m}} \approx 21 \text{ km/s} \left(\frac{T_0}{10^4 \text{ K}} \right)^{1/2}$$

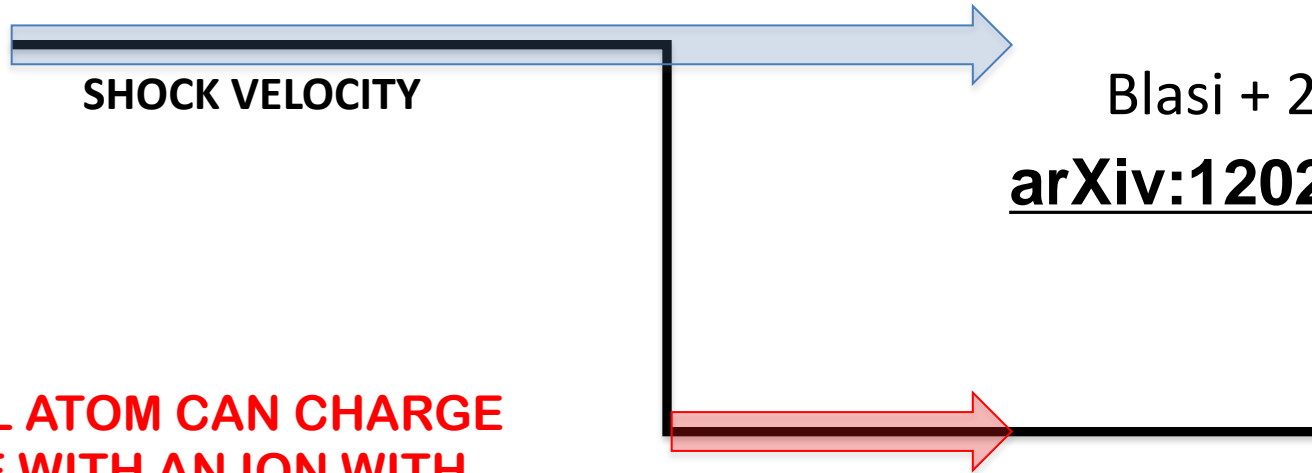


$$W_n \sim 30 - 50 \text{ km/s} \rightarrow T \sim 2 - 6 \cdot 10^4 \text{ K}$$

NARROW BALMER LINE BROADER THAN FOR AN UNMODIFIED SHOCK

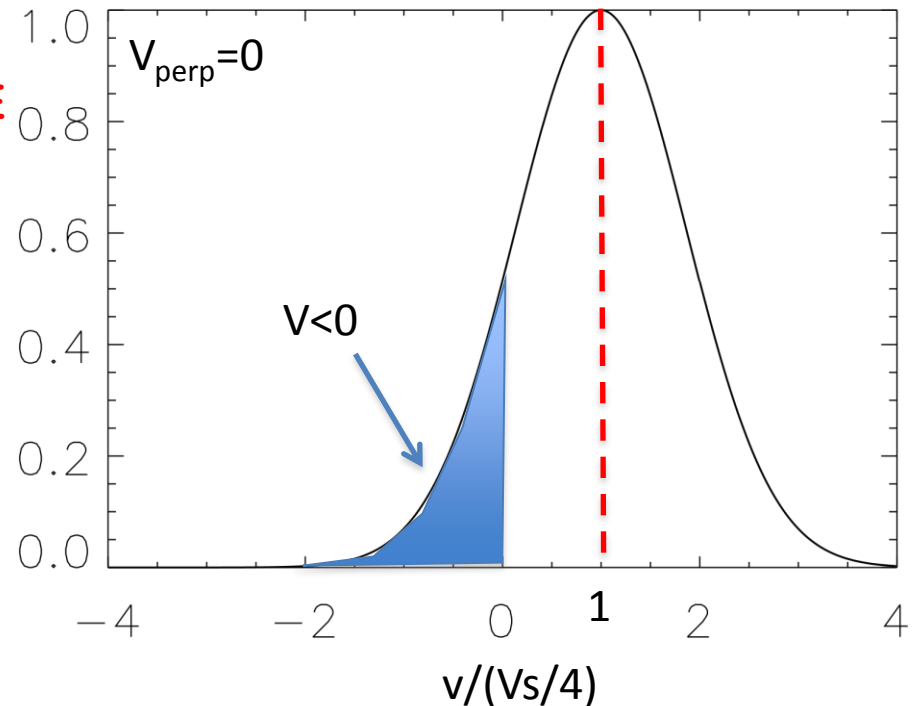
The neutral return flux

NEUTRALS
AND IONS

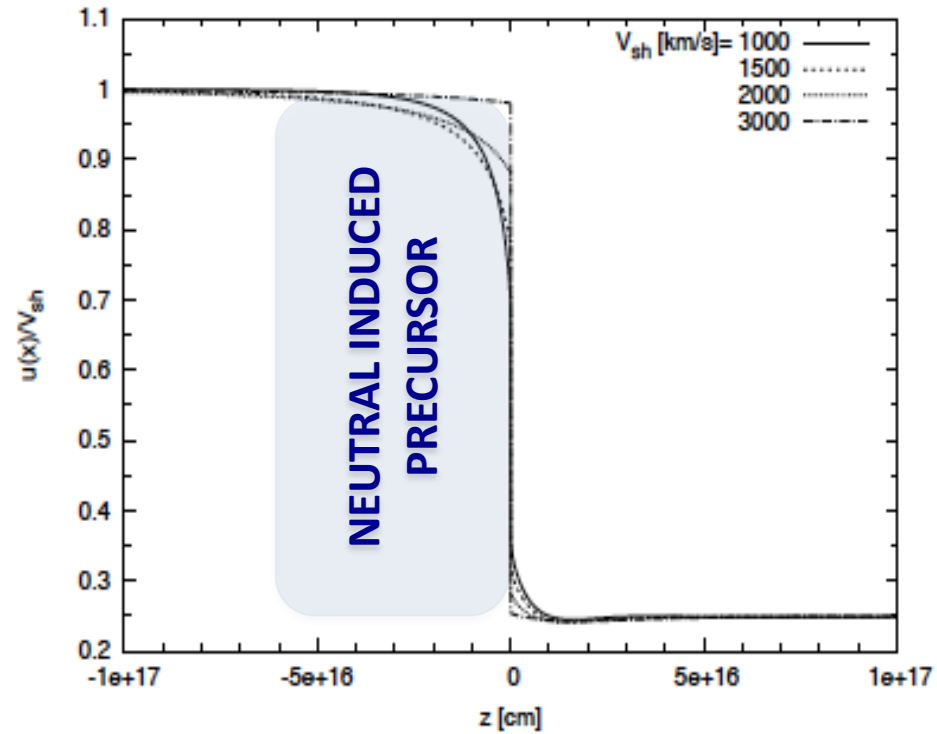
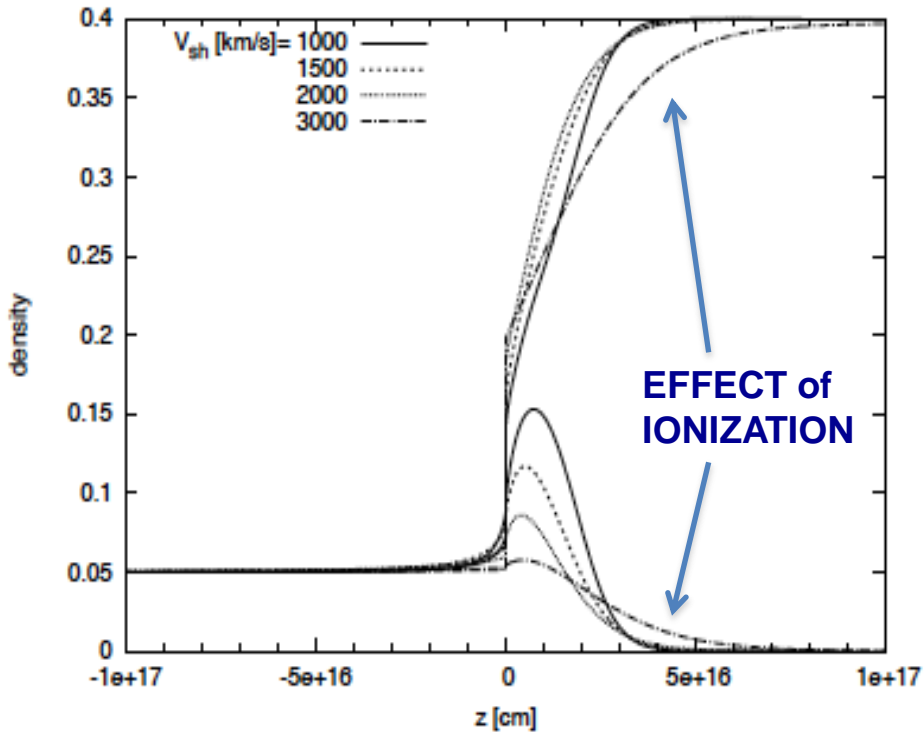


A NEUTRAL ATOM CAN CHARGE EXCHANGE WITH AN ION WITH $V < 0$, THEREBY GIVING RISE TO A NEUTRAL WHICH IS NOW FREE TO RETURN UPSTREAM

THIS NEUTRAL RETURN FLUX LEADS TO ENERGY AND MOMENTUM DEPOSITION UPSTREAM OF THE SHOCK!



SHOCK MODIFIED BY THE NEUTRAL RETURN FLUX

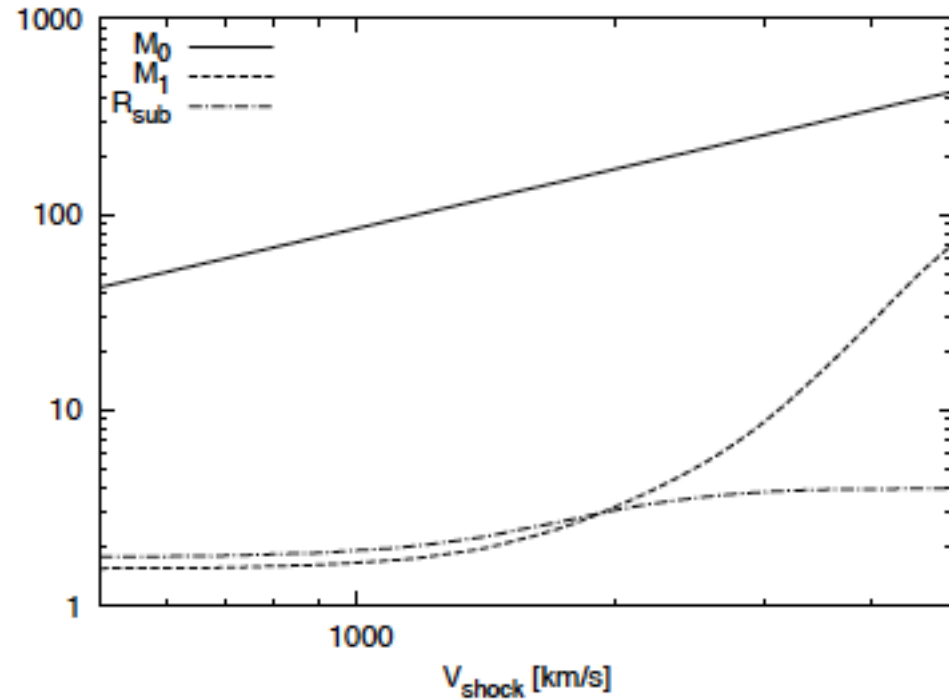
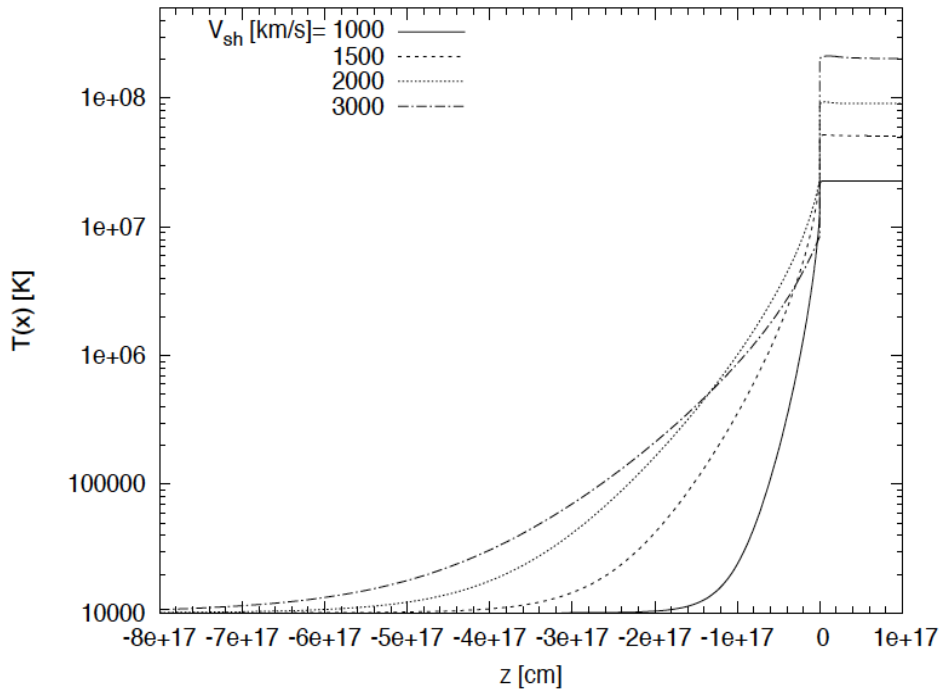


PB+2012

[arXiv:1202.3080](https://arxiv.org/abs/1202.3080)

NEUTRAL INDUCED PRECURSOR

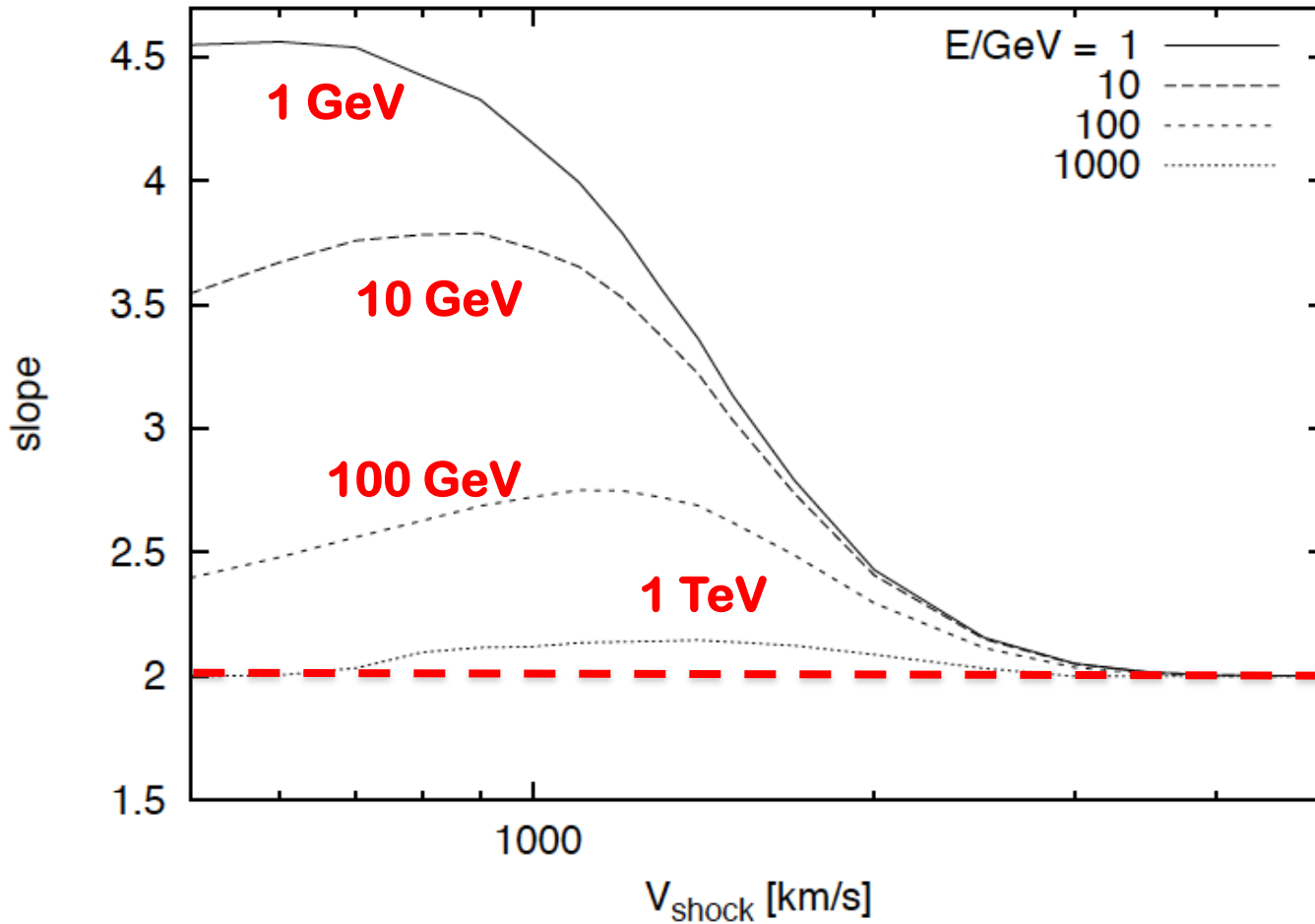
PB+2012 [arXiv:1202.3080](https://arxiv.org/abs/1202.3080)



EVEN FOR A STRONG SHOCK ($M \gg 1$) THE EFFECTIVE MACH NUMBER OF THE PLASMA IS DRAMATICALLY REDUCED DUE TO THE ACTION OF THE NEUTRAL RETURN FLUX

ACCELERATION OF TEST PARTICLES

PB+ 2012 [arXiv:1202.3080](https://arxiv.org/abs/1202.3080)



CONCLUSIONS

THE SN PARADIGM EXPLAINS THE GENERAL OBSERVATIONAL PICTURE OF THE ORIGIN OF COSMIC RAYS

BUT MISSING PIECES: STEEP SPECTRA, ANISOTROPY, SPECTRAL FEATURES IN THE NUCLEAR SPECTRA

MAGNETIC FIELD AMPLIFICATION INDUCED BY CR IS NOW ONE OF THE HOT TOPICS IN EXPLAINING X-RAY MORPHOLOGY AND REACH THE KNEE

THE KNEE CAN BE EXPLAINED REASONABLY WELL IN TERMS OF CHANGE OF CHEMICAL COMPOSITION, BUT TRANSITION WELL BELOW THE ANKLE

PRESENCE OF NEUTRALS EXTREMELY IMPORTANT DIAGNOSTIC TOOL FOR CR ACCELERATION AS WELL AS POSSIBLE CAUSE FOR SPECTRAL STEEPENING